

Upper critical field (H_{c2}) scaling near a quantum critical point in the heavy-fermion compound CeRhIn_5

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Abstract

We have measured the specific heat of CeRhIn_5 at extreme conditions of low temperatures (T), high pressures (P), and high magnetic fields ($H \perp c$ -axis). Discrete forms of the upper critical field H_{c2} scaling are observed below and above the critical pressure $P_{c1}(= 1.75 \text{ GPa})$. For $P < 1.75 \text{ GPa}$, where unconventional superconductivity and magnetism coexist, H_{c2} linearly increases with decreasing temperature. For $P > 2.35 \text{ GPa}$, where only unconventional superconductivity is observed up to the normal state, H_{c2} shows normal behavior, saturating at low temperatures. In the intermediate pressure, $1.75 \text{ GPa} < P < 2.35 \text{ GPa}$, where applied magnetic field induces magnetism in the mixed superconducting state, H_{c2} shows a crossover from normal temperature dependence to the unusual linear T dependence. These observations delineate the interplay between superconductivity and antiferromagnetic fluctuations that exist in the vicinity of the quantum critical point.

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The heavy-fermion compound CeRhIn_5 is a model unconventional superconductor that reveals superconductivity when long-range antiferromagnetic (AFM) order is suppressed under pressure. Recently, a specific heat study revealed magnetism in the mixed superconducting state of CeRhIn_5 , suggesting a common relationship among hidden magnetism, quantum criticality, and unconventional superconductivity in cuprate and heavy-electron systems [1]. The details of how magnetic fluctuations interplay with superconductivity, however, have been rarely studied. Here, we report disparate forms of the upper critical field scaling below and above a critical pressure P_{c1} , where the AFM order is suppressed to zero temperature in the pressure-tuned heavy-fermion superconductor CeRhIn_5 , reflecting the interplay between superconductivity and antiferromagnetic fluctuations near the critical point P_{c1} .

The temperature–pressure (T – P) phase diagram of CeRhIn_5 is shown in Fig. 1, which is typical of unconventional superconductors, including heavy fermion com-

pounds, organics, and high- T_c cuprates: in high- T_c cuprates, the control parameter (x -axis) is chemical substitution. With increasing pressure, the magnetic transition temperature (T_N) due to the incommensurate antiferromagnetic state with $\mathbf{Q} = (0.5, 0.5, 0.297)$ goes through a maximum at 0.55 GPa and then decreases. With further increasing pressure, superconductivity is induced within the antiferromagnetic state and coexists microscopically [2]. At P_{c1} , where T_N becomes equal to T_c , the long-range AFM order sharply disappears. This first-order or weakly first-order transition at P_{c1} is puzzling in that it does not ensure its clear connection to strange metallic behavior or non-Fermi liquid behavior observed in the normal state above the superconducting dome. Recently, it was found that the magnetism is hidden by unconventional superconductivity and can be revealed by applying a magnetic field, resolving the dichotomy because it now provides a smooth evolution of T_N into the superconducting dome, terminating at P_{c2} , an antiferromagnetic quantum critical point [1]. Our phase diagram determined by specific heat measurements reveals several crucial aspects that have not been appreciated previously. First, the superconducting

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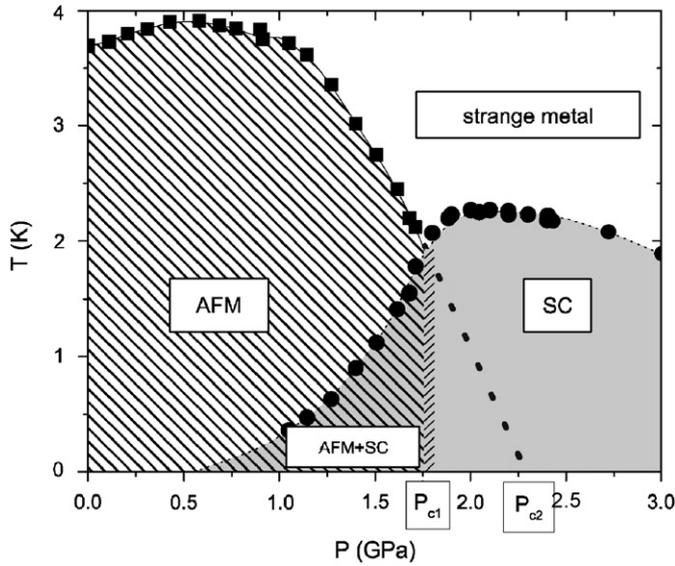


Fig. 1. Temperature–pressure (T – P) phase diagram of CeRhIn_5 determined from specific heat measurements at $H = 0$ kOe. AFM and SC stand for antiferromagnetic and superconducting states, respectively. P_{c1} and P_{c2} are explained in the text.

phase in the AFM state now extends well below the critical pressure P_{c1} , delineating that the SC phase is a bulk property, not a filamentary phase as has been speculated before [3]. The experimentally determined T_c is 0.36 K at 1.05 GPa and can be smoothly extended to 0 K at 0.55 GPa, where T_N is of maximal value. Second, the superconducting dome is asymmetric: T_c changes slowly with pressure for $P > P_{c1}$, while T_c is dramatically suppressed with decreasing pressure for $P < P_{c1}$, indicating an interplay between magnetism and superconductivity. The sharp suppression of T_c in the coexisting phase has been also reported in CeCu_2Si_2 , another family of heavy fermion superconductors [4].

The interplay between magnetism and unconventional superconductivity is also manifested in the upper critical field of CeRhIn_5 . Fig. 2a representatively shows H_{c2} as a function of reduced temperature $t = T/T_{c0}$, where T_{c0} is the SC transition temperature at zero field and is 1.41, 2.27, and 2.18 K for 1.62, 2.1, and 2.45 GPa, respectively. For $P > P_{c2}$, which is represented by $P = 2.45$ GPa (circles), the temperature dependence of H_{c2} is similar to that of CeCoIn_5 , where $H_{c2}(0)$ is strongly limited due to Pauli limiting effects. The Maki parameter $\alpha (= H_{c2}^0/H_P)$ of CeRhIn_5 , which is a measure of the relative strength of the orbital limiting field H_{c2}^0 and the Pauli field H_P , is 6.3 ± 0.8 in the weak coupling limit. Considering that CeRhIn_5 is in the very clean limit with mean free path of order microns, this sufficiently large α (for $\alpha > 1.8$) suggests this pressure-induced superconductor as an ideal candidate to explore the inhomogeneous superconducting phase or Fulde and Ferrell and Larkin and Ovchinnikov (FFLO) state [5]. At $P = 1.62$ GPa (triangles), where superconductivity coexists with the incommensurate AFM state, H_{c2} becomes surprisingly linear down to lowest temperature ($t = 0.2$).

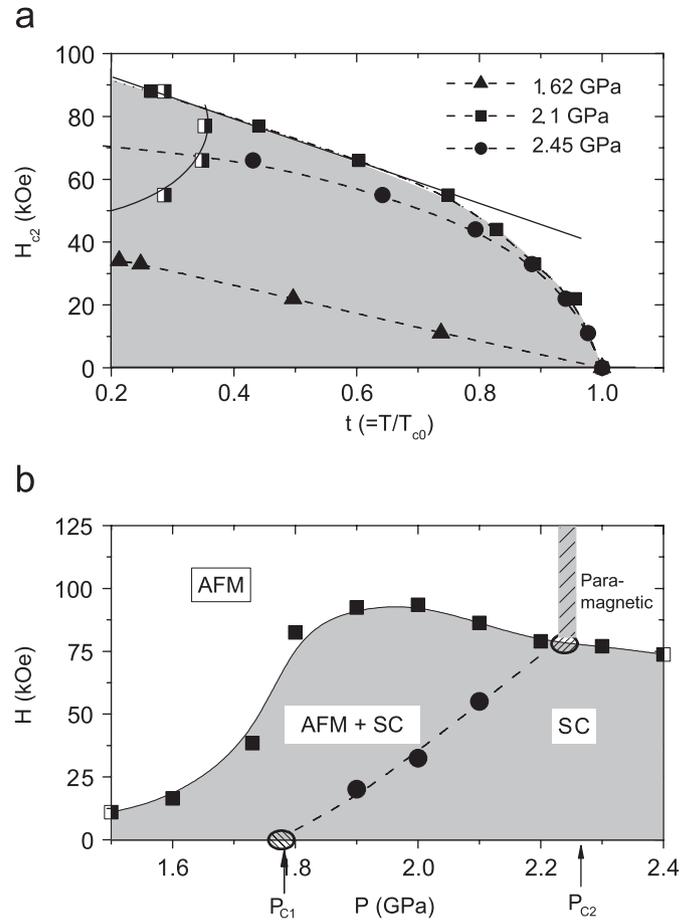


Fig. 2. (a) Upper critical field (H_{c2}) versus reduced temperature T/T_{c0} at 1.62 (triangles), 2.1 (squares), and 2.45 GPa (circles), where T_{c0} is SC transition temperature at zero field and is 1.41, 2.27, and 2.18 K, respectively. The half-filled squares describe field-induced magnetism at 2.1 GPa. Magnetic field is applied parallel to ab-plane. Lines are guides to eyes. (b) H – P phase diagram at 650 mK, which represents a $T = 0$ K quantum phase transition [1]. Squares and circles describe the upper critical field H_{c2} and field-induced transition H_M from a pure SC state to the coexisting phase of SC and AFM states, respectively. The two critical pressures P_{c1} and P_{c2} are explained in the text.

Recently, a linear- T dependence of H_{c2} is also reported in the pyrochlore superconductor KOs_2O_6 , which is argued to be due to band effects [6]. In CeRhIn_5 , the pure band effects, however, may be less relevant than in KOs_2O_6 . Above the critical pressure P_{c1} (squares in Fig. 2a), where only superconductivity is observed for $H = 0$, the H_{c2} shows normal behavior near T_c , similar to high-pressure data for $P > P_{c2}$. At this pressure, applied magnetic field induces long-range magnetic order in the superconducting state (see Fig. 2b) and the H_{c2} temperature dependence changes from normal to linear behavior near a field value at which field-induced AFM appears, suggesting that the AFM order is crucial to the anomalous T -linear H_{c2} .

We have reported a comprehensive T – P phase diagram of the heavy-fermion compound CeRhIn_5 , which shows bulk superconductivity in the AFM state well below P_{c1} . The strong asymmetry of the superconducting dome below

and above P_{c1} and the disparate forms of upper critical field scaling strongly suggest the interplay between antiferromagnetic fluctuations and unconventional superconductivity. To finish this brief report, it seems fitting to note a theoretical work that has raised the possibility of exotic superconductivity, such as odd-frequency p-wave spin singlet state in the vicinity of a quantum critical point [7].

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